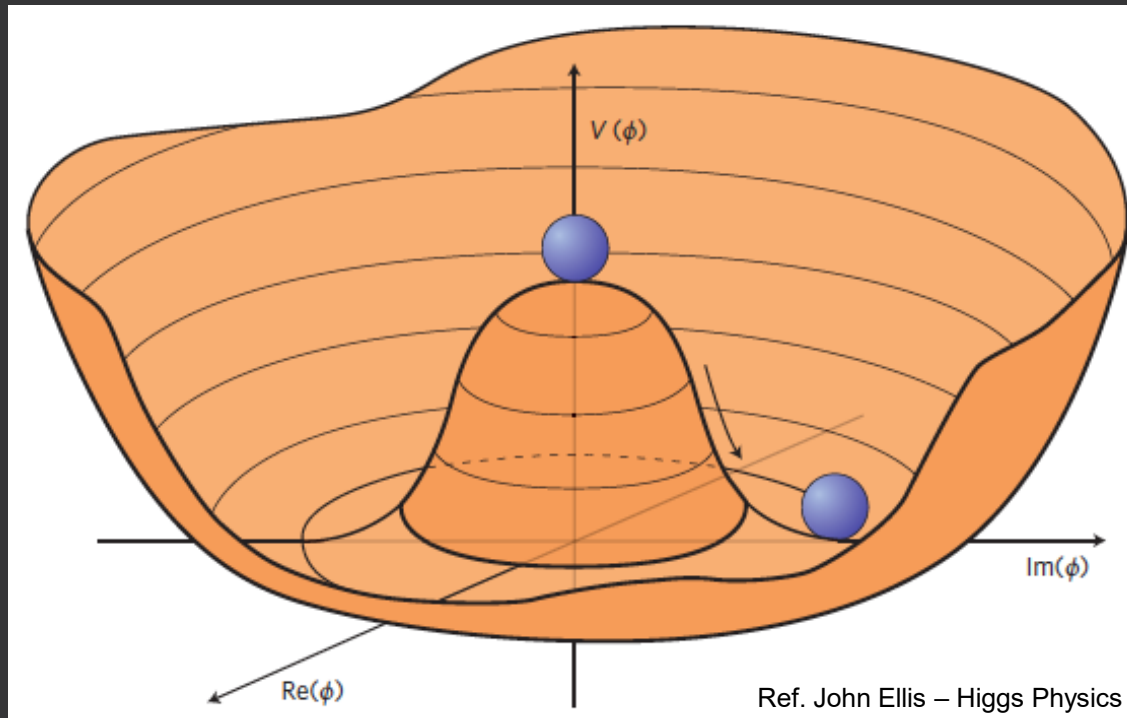


The Higgs as the Inflaton

Patrick Adams, Advisor: Daniel Grin

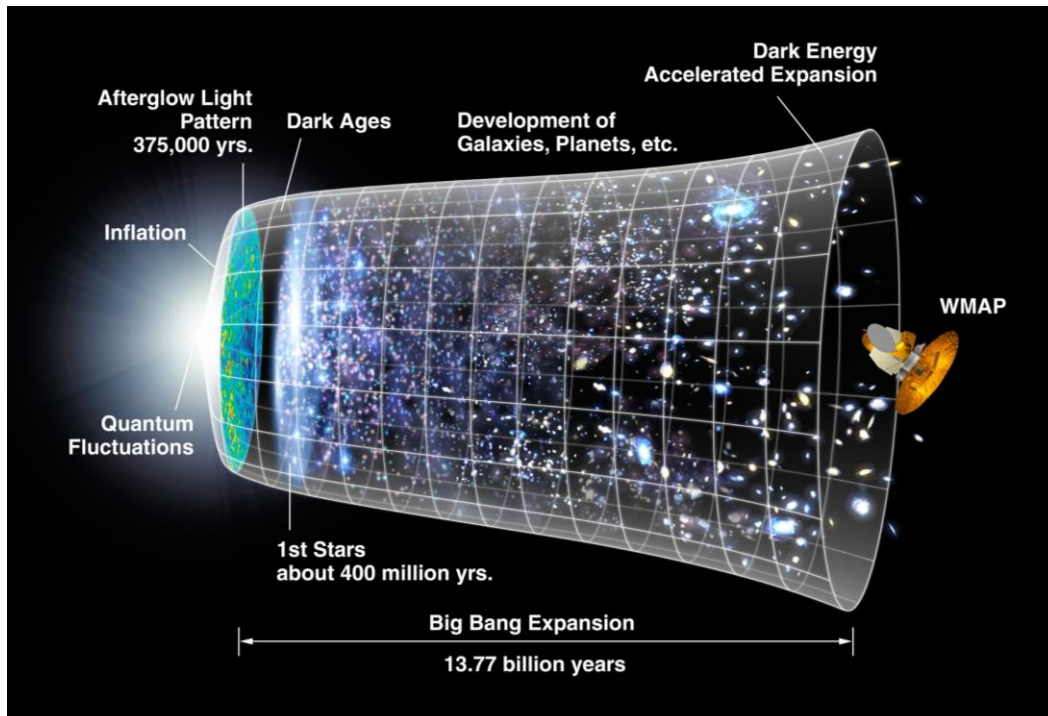


Our Collaborators

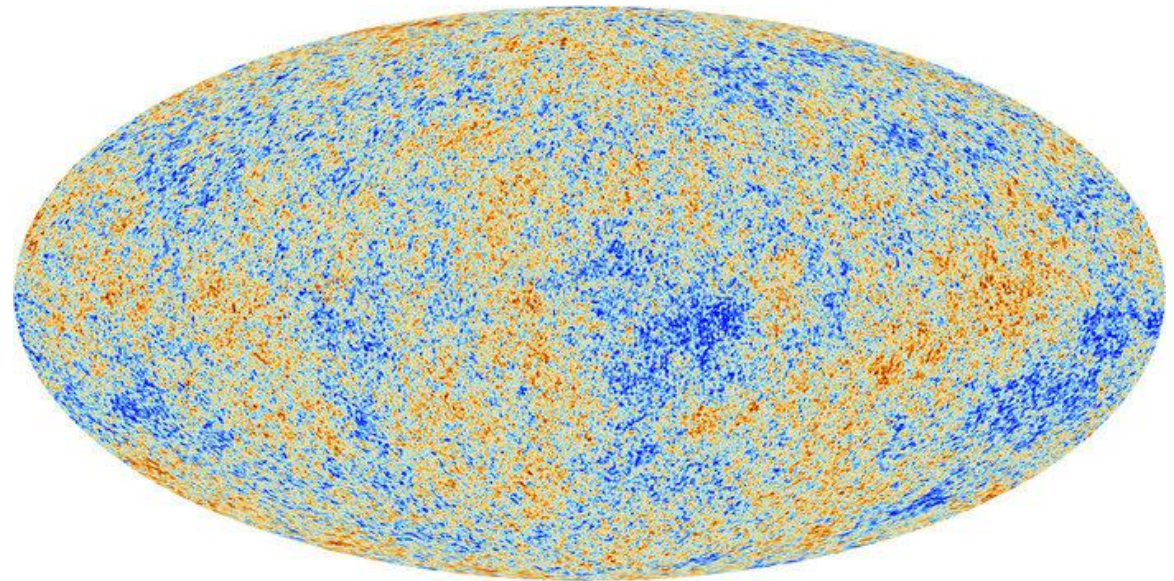
- **David Kaiser** – MIT – Wrote detailed notes on Higgs inflation
- **Renée Hložek** – University of Toronto – Project used Fisher analysis to test how compatible a simplified New Higgs inflation is with the CMB
- **William Koster** – Haverford Alum – Thesis work with Prof. Hložek inspired this project
- **Chanda Prescod-Weinstein** – Brown University

Steady State to Big Bang

- Detection of the Cosmic Microwave Background (CMB) suggests there was a Big Bang
- The universe expanded from a singularity
- Expanded from more dense state to less dense state



Ref. NASA/WMAP Science Team

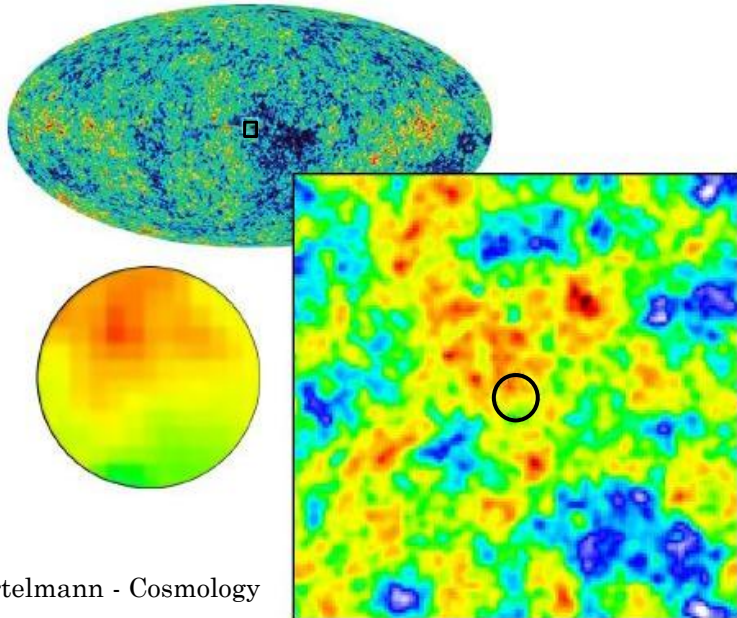


Ref. ESA and the Planck Collaboration

Problems with the Big Bang

Horizon Problem

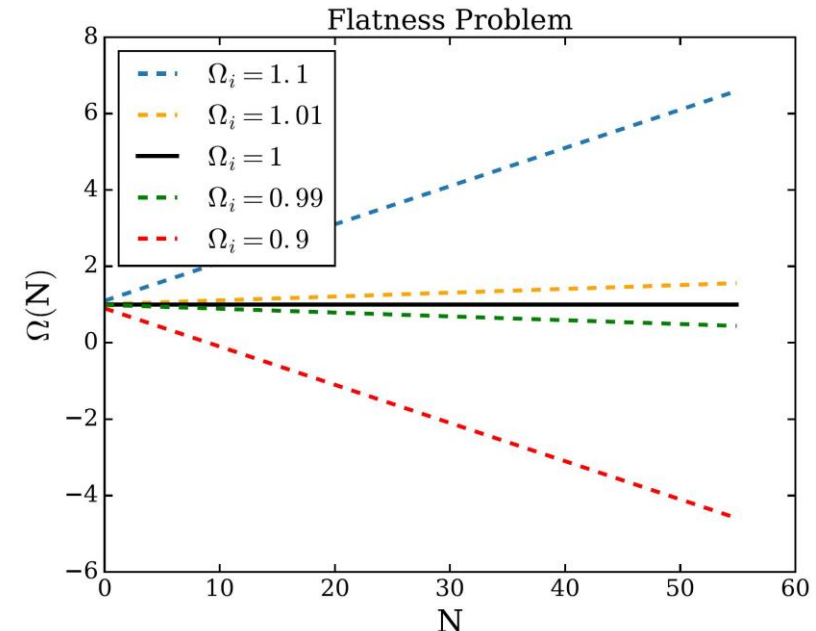
- CMB is isotropic and homogenous: about the same temperature
- The “particle horizon” is only about 1° , but cosmic horizon must have been much larger



Ref. Bartelmann - Cosmology

Flatness Problem

- Flat universe today: More flat in past
- In a matter and radiation dominated universe, perturbation size would only increase

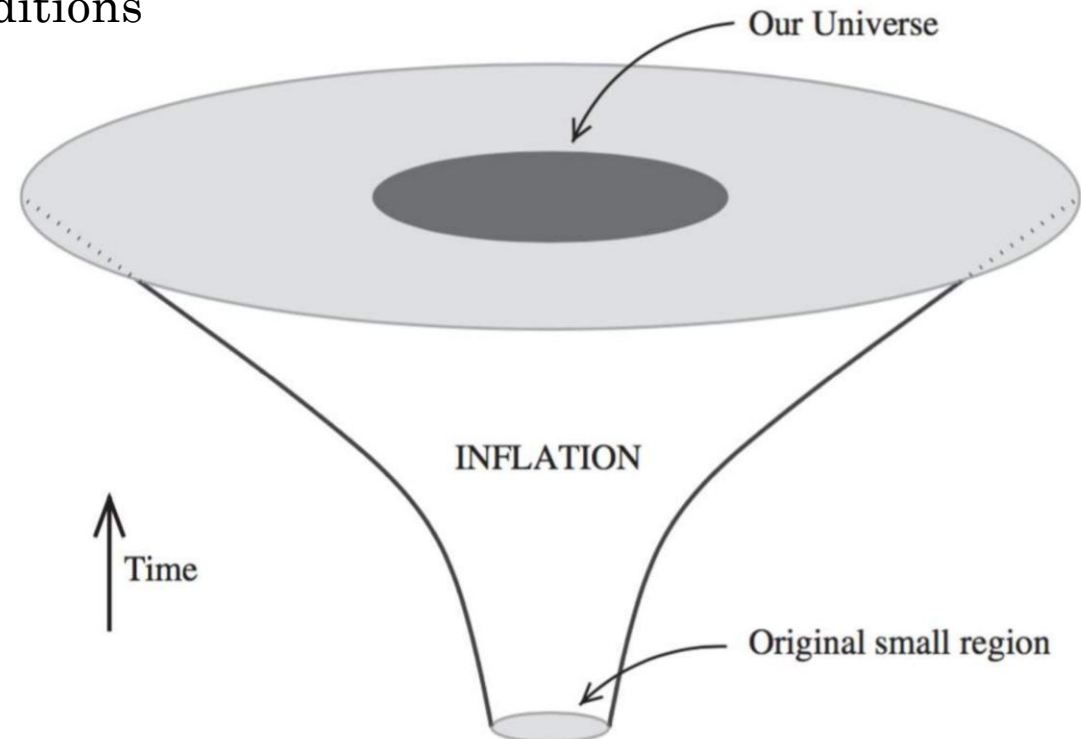


Inflation as a Solution

- Epoch of accelerated expansion of the universe
- Inflation allows for a period in which the universe is in causal contact
- Avoids need for very specific initial conditions
- First Proposed by Prof. Alan Guth

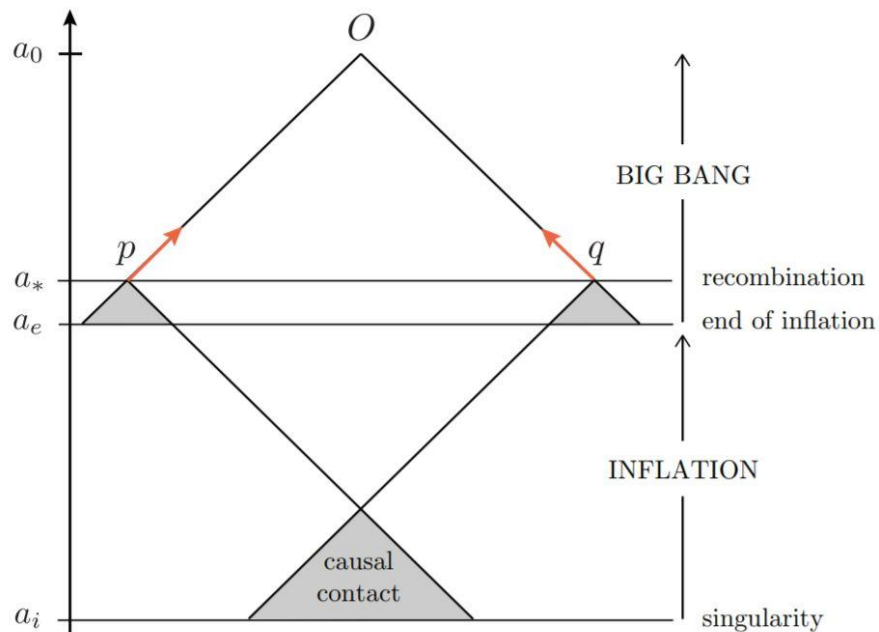
N is E-Folding Value:

$$a(t) = e^{\int_{t_i}^t H dt'}$$
$$= e^N$$



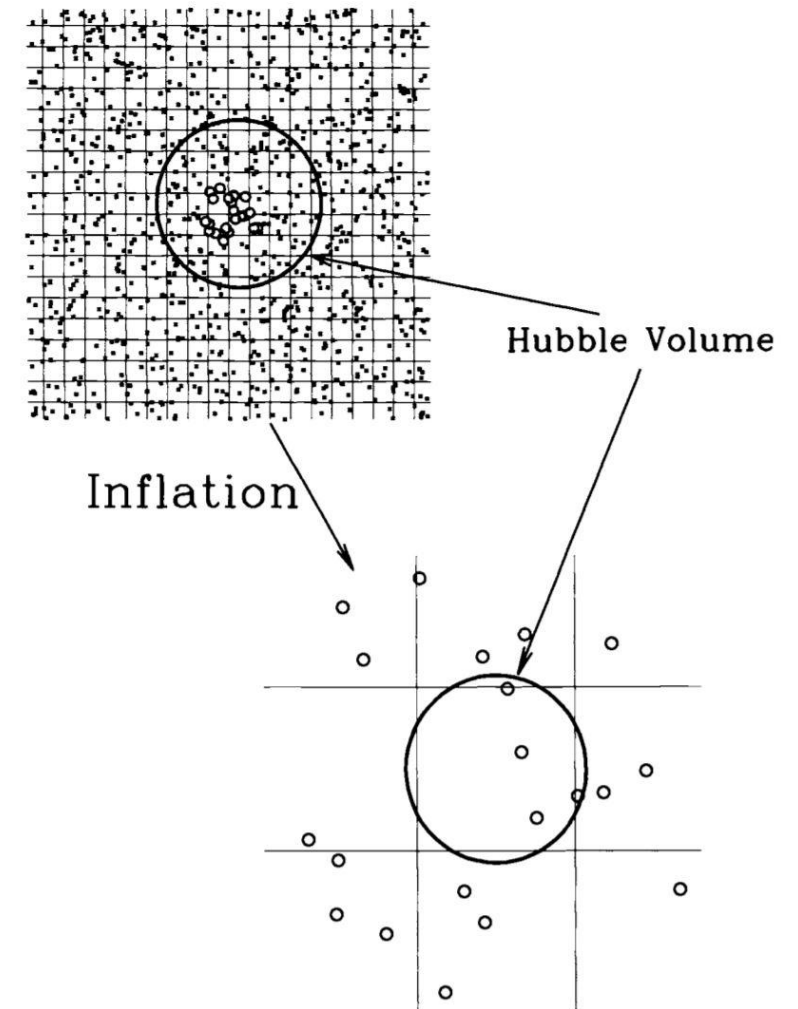
“Shrinking” Causal Horizon

- Universe's co-moving coordinates expanded during inflation – effective causal horizon stays the same
- Thus, the effective causal horizon that matter can interact with shrinks



Ref. Baumann – TASI Lectures on Inflation

Ref. Dodelson – Modern Cosmology

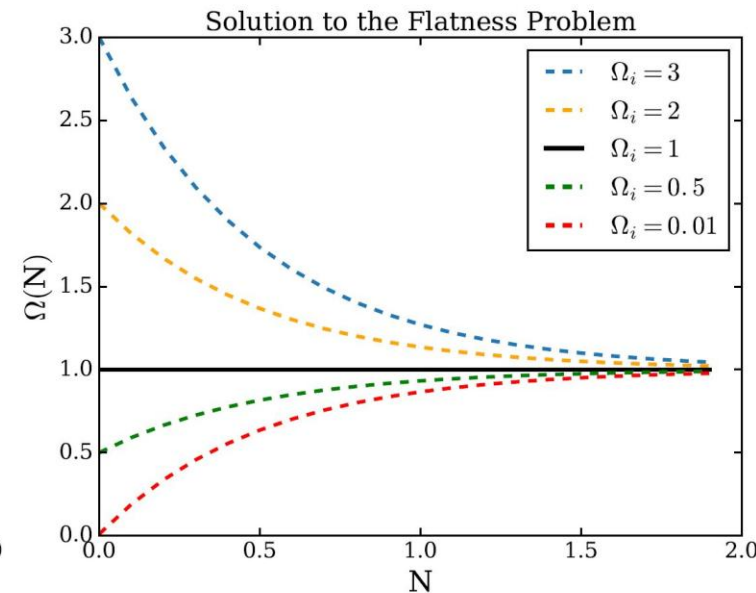
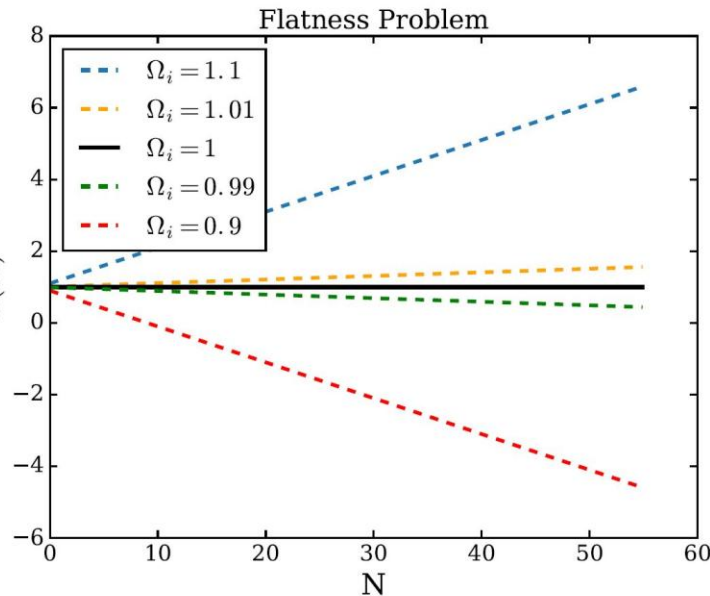


Inflating Universe to Flat Universe

- Density perturbations change exponentially during inflation, thus initial conditions converge to a flat universe
- A radiation or matter dominated universe tends to increase the relative size of perturbations

$$|1 - \Omega(t)| = \frac{k}{a(t)^2 H(t)^2}$$

$$|1 - \Omega(t)| \propto t^{\Omega_i}$$



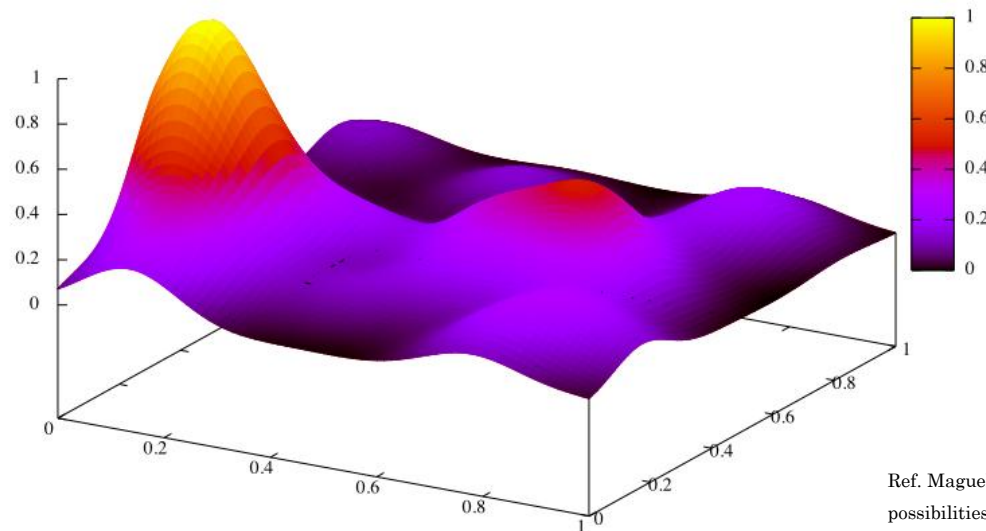
$$|1 - \Omega(t)| \propto e^{Ht}$$

Inflation described by a scalar field

- Homogenous field plus a perturbed field:

$$\phi(t, \vec{x}) = \varphi(t) + \delta\phi(t, \vec{x})$$

- Associated spin-0 particle: the *Inflaton*
- Ex. Gravitational Potential, Temperature

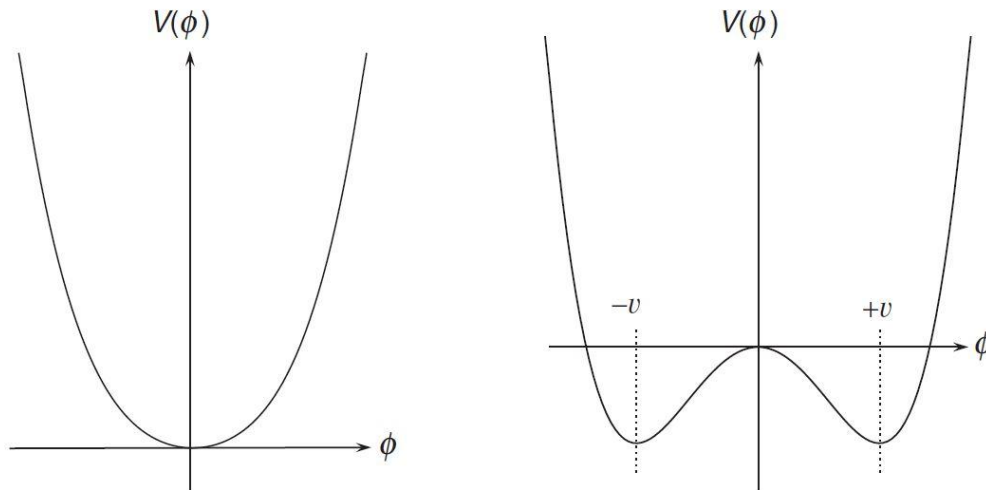


Ref. Maguel Negrão - The challenges and possibilities of real-time Wave Field Synthesis

Higgs Boson as the *Inflaton*

- Higgs Boson is only known scalar field and spin-0 particle
- Discovered in 2012 at CERN, questions linger about whether it is a candidate for inflation
- Shape of Higgs Potential is ideal for inflation

$$V(\phi) = \lambda(\phi^2 - M^2)^2$$



Ref. Thomson – Modern Particle Physics

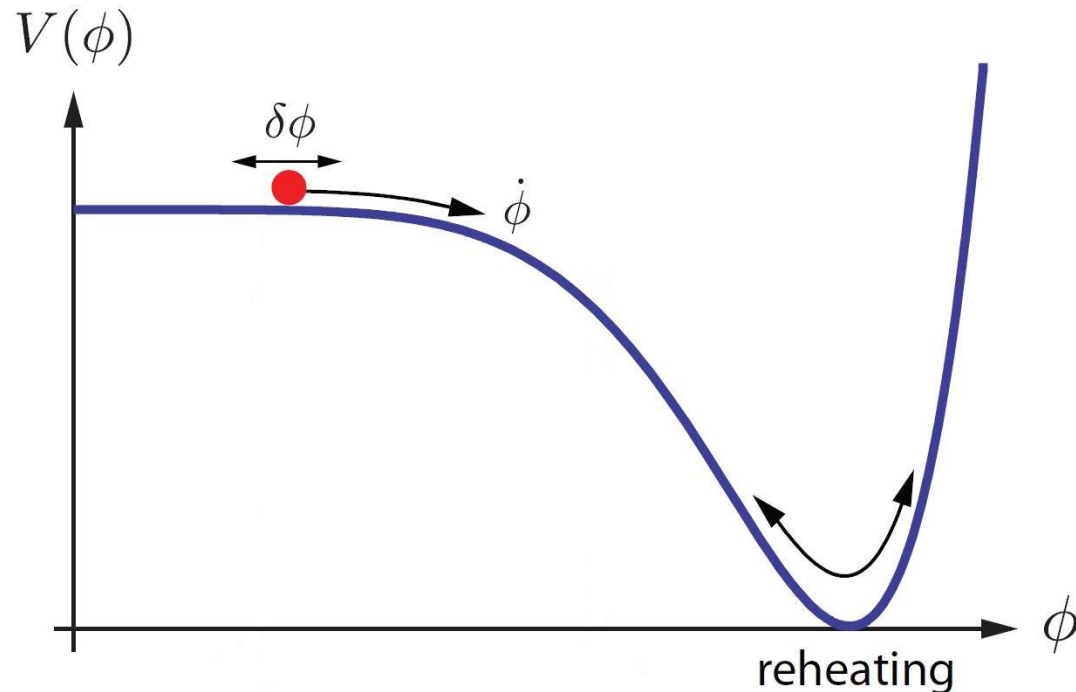
Standard Model of Elementary Particles

three generations of matter (fermions)			interactions / force carriers (bosons)
I	II	III	
mass $\approx 2.2 \text{ MeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ up	mass $\approx 1.28 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ charm	mass $\approx 173.1 \text{ GeV}/c^2$ charge $\frac{2}{3}$ spin $\frac{1}{2}$ top	mass $\approx 125.09 \text{ GeV}/c^2$ charge 0 spin 0 higgs
 down	 strange	 bottom	 gluon
 electron	 muon	 tau	 photon
 electron neutrino	 muon neutrino	 tau neutrino	 Z boson
			 W boson

Ref. Wikimedia Commons

Expansion as a “slow-roll”

- Inflation’s expansion is controlled by the shape of field’s potential
- Slow-roll parameters yield dynamics



Slow-Roll Conditions

$$|\epsilon(\phi)| \ll 1$$

$$|\eta(\phi)| \ll 1$$

Slow-Roll Parameters

$$\epsilon(\phi) = \frac{M_{\text{pl}}^2}{2} \left(\frac{1}{V(\phi)} \frac{dV(\phi)}{d\phi} \right)^2$$

$$\eta(\phi) = M_{\text{pl}}^2 \left(\frac{1}{V(\phi)} \frac{d^2V(\phi)}{d\phi^2} \right)$$

Inflation's end

- Dictated by the Friedman Equation and equation of state parameter w :

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3c^4}(\rho + 3p) \quad w = \frac{p}{\rho} = \frac{\frac{1}{2}\dot{\phi}^2 - V(\phi)}{\frac{1}{2}\dot{\phi}^2 + V(\phi)}$$

- Inflation occurs when pressure dominates over matter density

$$3p < -\rho \Rightarrow w < -\frac{1}{3}$$

Equations of Motion

- **“Old Higgs” model** for inflation – assumes GR holds for all time

$$S = \int d^4x \sqrt{-g} \left(\frac{M_{\text{pl}}^2}{2} R - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right)$$

- Extremizing the act yields the homogeneous Equations of Motion

$$\ddot{\phi} + 3H\dot{\phi} + \frac{\partial V(\phi)}{\partial \phi} = 0$$

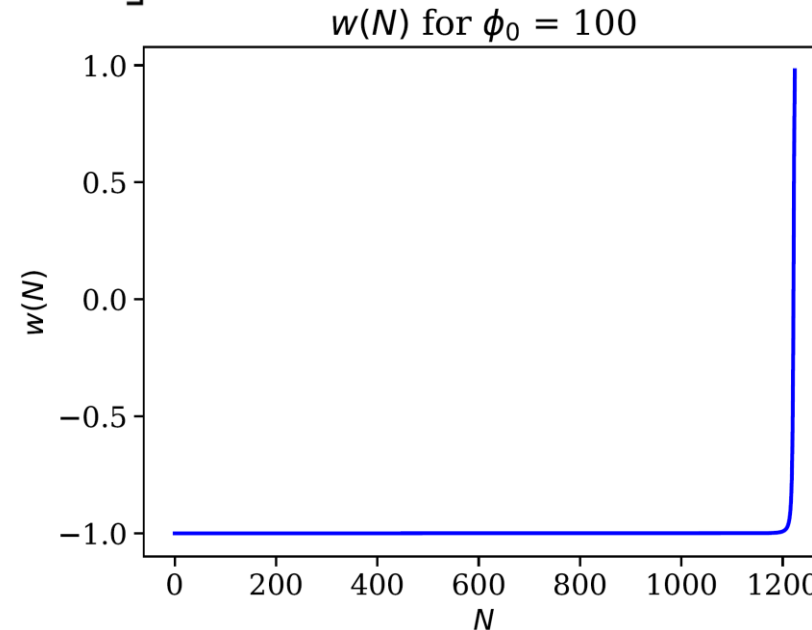
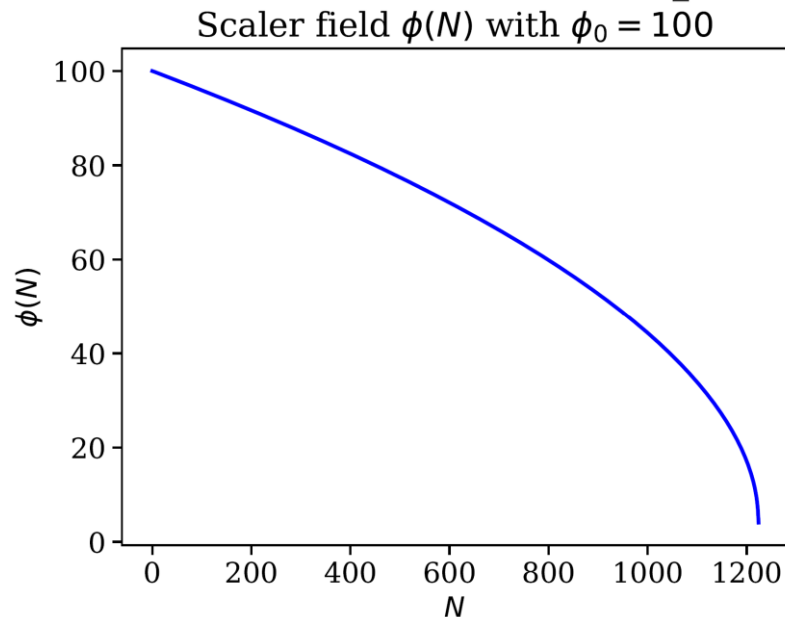
$$H^2 = \frac{8\pi G}{3} \left[\frac{1}{2} \dot{\phi}^2 + V(\phi) \right]$$

Rewriting Equations to see Inflation

- Rewrote equations in terms of E-folding parameter N for Higgs potential

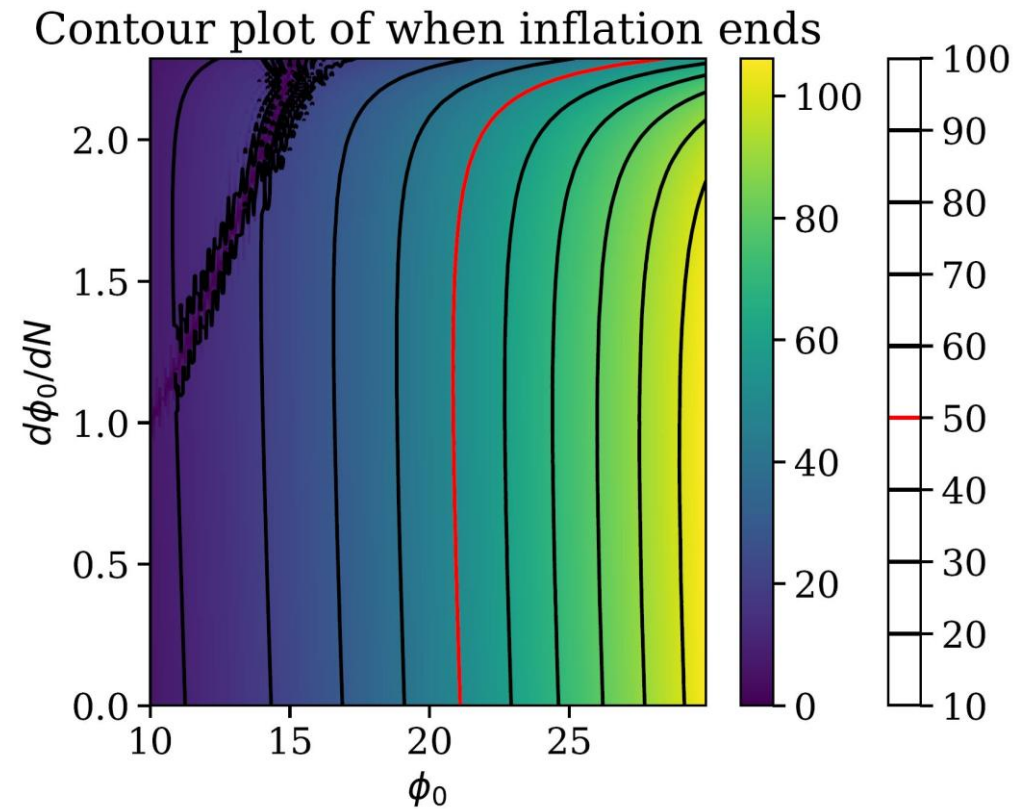
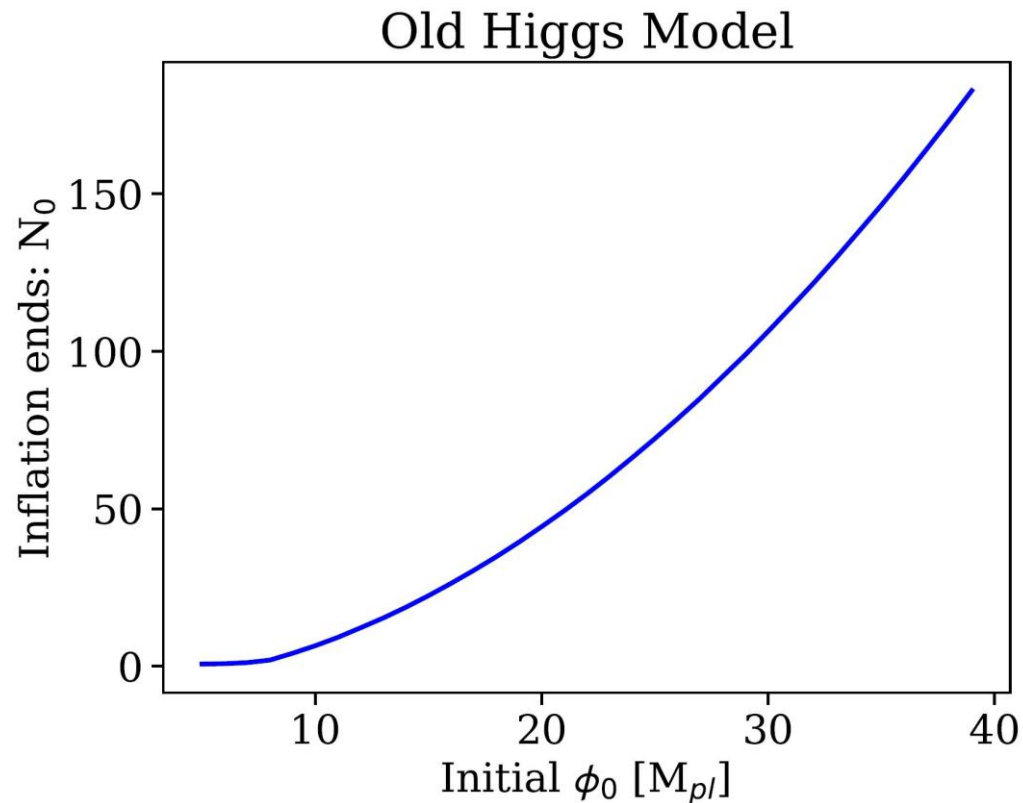
$$\tilde{H}^2 \frac{d^2 \tilde{\varphi}}{dN^2} + (3\tilde{H} + \frac{d\tilde{H}}{dN}) \tilde{H} \frac{d\tilde{\varphi}}{dN} + 4\lambda(\tilde{\varphi}^3 - \frac{M^2}{M_{\text{pl}}^2} \tilde{\varphi}) = 0$$

$$\tilde{H} = \left[\frac{\lambda(\tilde{\varphi}^2 - \frac{M^2}{M_{\text{pl}}^2})^2}{3 - \frac{1}{2}(\frac{d\tilde{\varphi}}{dN})^2} \right]^{1/2}$$



Old Higgs Phase Space

- Solved for inflation's end for different initial scalar field values
- Additionally varied initial $\frac{d\phi}{dN}$ versus ϕ_0 value



Alas, Old Higgs is no good

- Allows for inflating universe but is constrained by particle physics
- Higgs self interaction constant is $0.11 \leq \lambda < 0.27$

$$H^2 \approx \frac{\lambda}{12M_{\text{pl}}^2} \phi^4$$

$$\dot{\phi} \approx -\frac{\lambda}{3H} \phi^3$$

- Applying a quantum gravity and slow-roll constraint yields:

$$\lambda \ll 10^{-2}$$

Non-minimal coupling

- “**New Higgs**” model’s couples its kinetic energy term directly to the Einstein curvature tensor

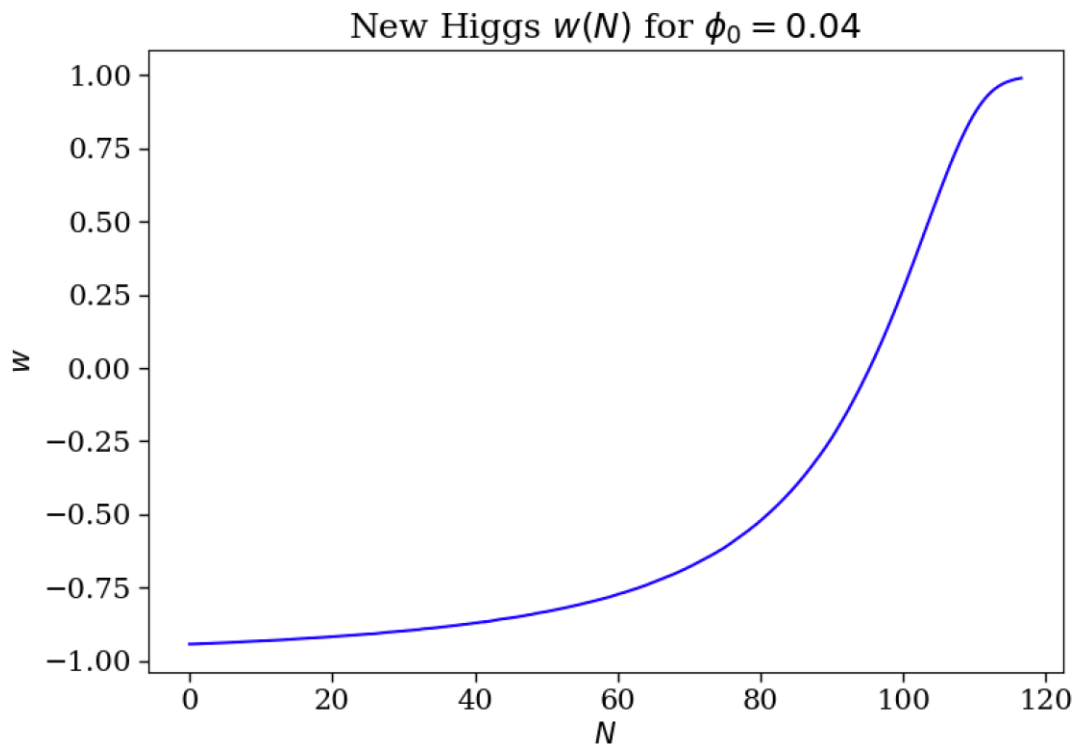
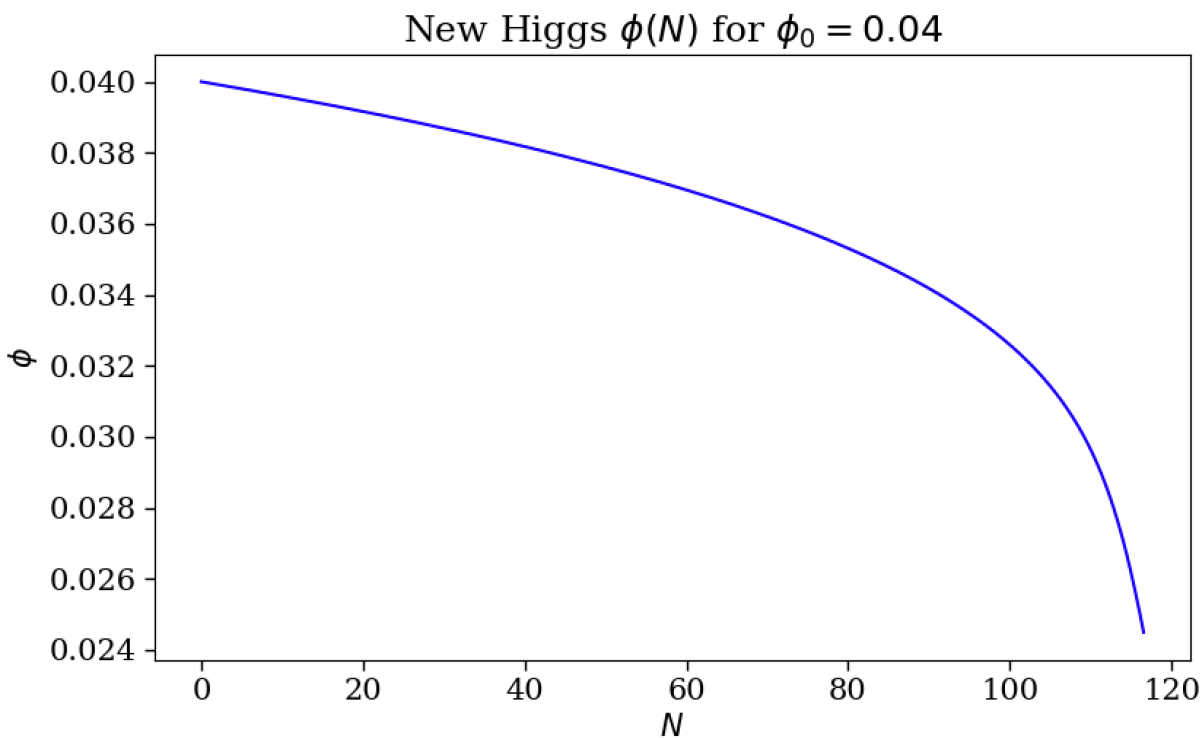
$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} R - \frac{1}{2} \left(g^{\mu\nu} - \frac{1}{\mu^2} G^{\mu\nu} \right) \partial_\mu \phi \partial_\nu \phi - V(\phi) \right]$$

- In the limit of $\phi \gg \frac{1}{\mu^2}$, the EOM looks like the Old Higgs Equations

$$\ddot{\varphi} \left[1 + \frac{3H^2}{\mu^2} \right] + 3H\dot{\varphi} \left[1 + \frac{1}{\mu^2} \left(3H^2 + 2\dot{H} \right) \right] + \frac{\partial V(\varphi)}{\partial \varphi} = 0$$
$$H^2 = \frac{1}{3M_{\text{pl}}^2} \left[\frac{1}{2} \dot{\varphi}^2 \left(1 + \frac{9H^2}{\mu^2} \right) + V(\varphi) \right]$$

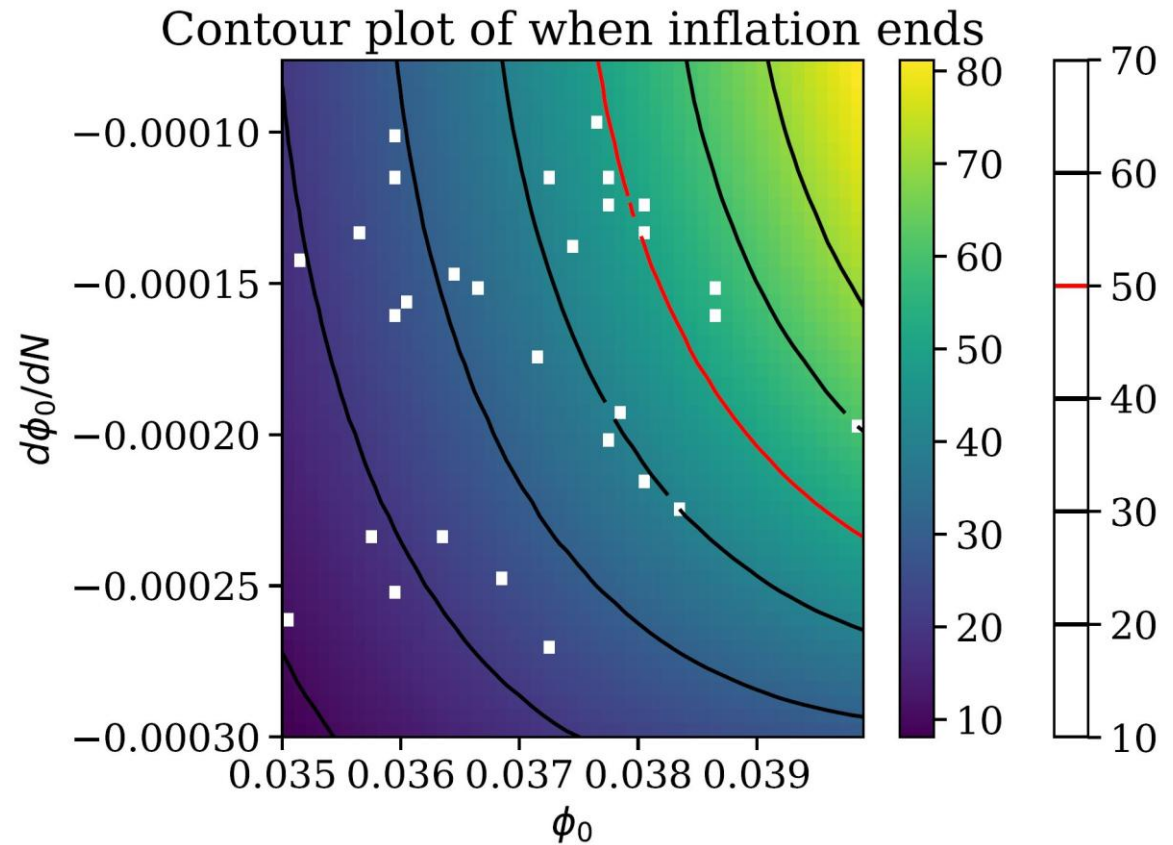
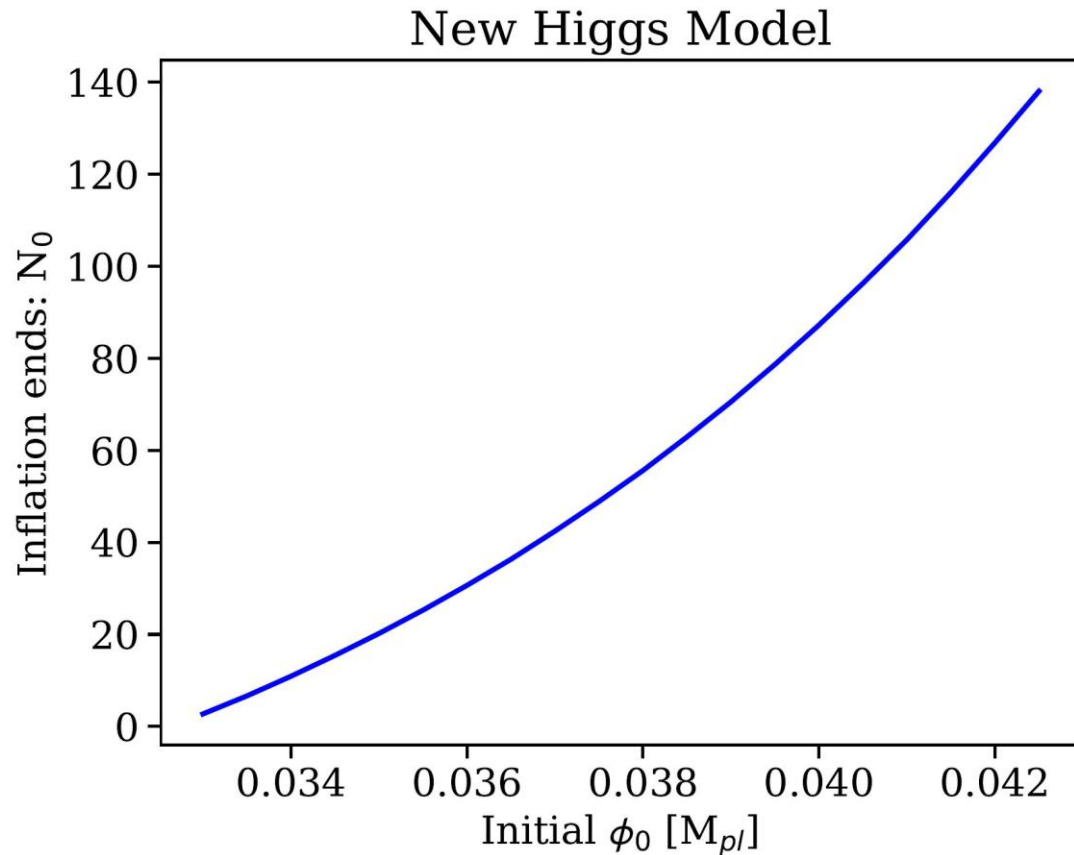
New Higgs Dynamics

- I solved for the New Higgs equations directly
- This required rewriting the equations in terms of e-foldings



New Higgs Phase Space

- Some convergence issues in the right figure, left figure is converged!



Two Candidates, Two Potentials

Old Higgs Inflation

- Not Coupled:

$$\Phi = \Psi$$

- Assumes GR holds for all time
- Ruled out as we have previously

New Higgs Inflation

- Non-Minimally Coupled:

$$\Phi \neq \Psi$$

- Assumes GR may not hold prior to CMB
- Allows for a shear between the space-like Φ and time-like Ψ potentials

FRW Metric which describes spacetime in cosmology:

$$ds^2 = a(\tau)^2 [-(1 + 2\Psi)d\tau^2 + (1 - 2\Phi)(dr^2 + r^2 d\Omega^2)]$$

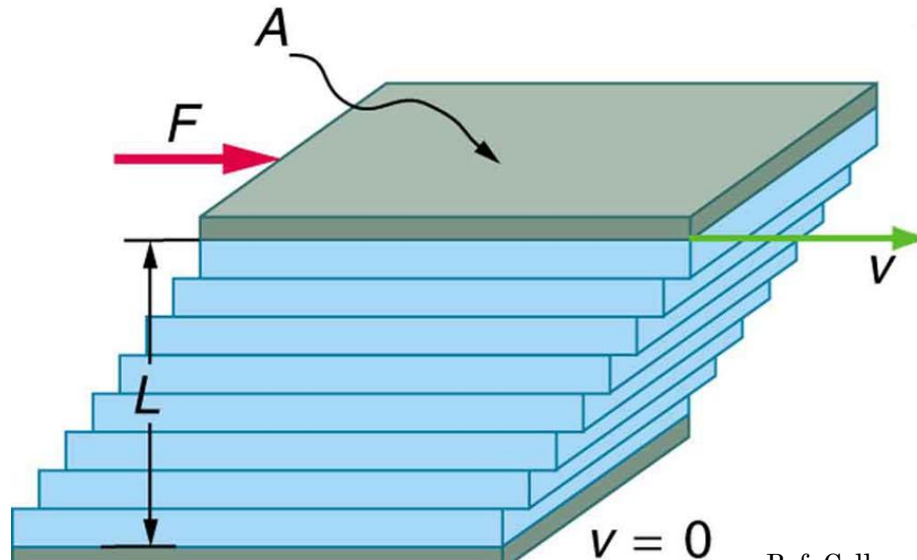
Early Universe Shear Stress

Fourier sum over matter species i :

$$k^2\Phi = -4\pi Ga^2 \sum_i \rho_i \Delta_i$$

$$k^2(\Psi - \Phi) = -12\pi Ga^2 \sum_i \rho_i (1 + w_i) \sigma_i$$

If shear $\sigma_i = 0$,
then $\Psi = \Phi$.

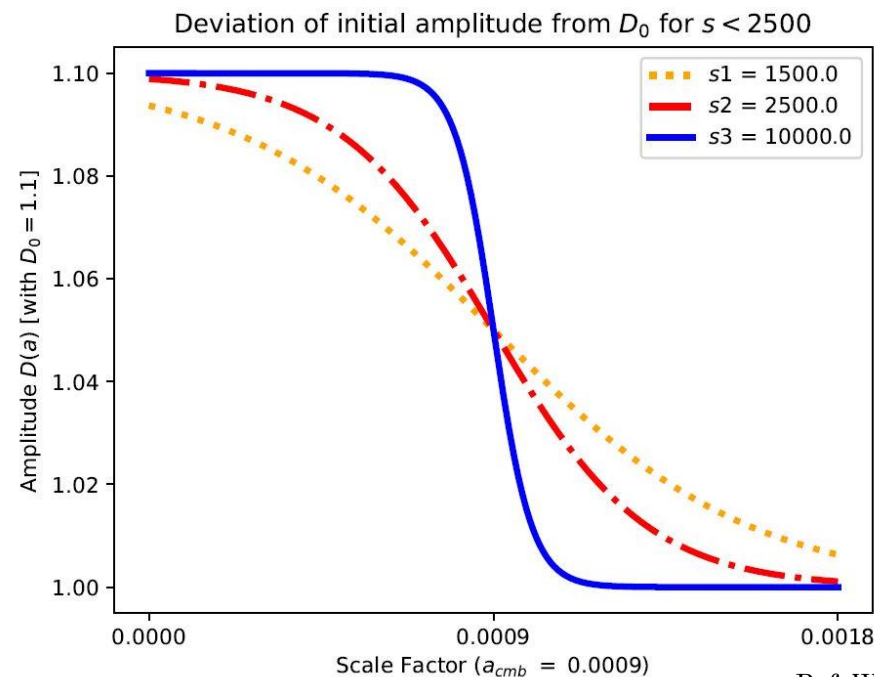


Previous New Higgs Research

- Work by William Koster tested a simplified model such that:

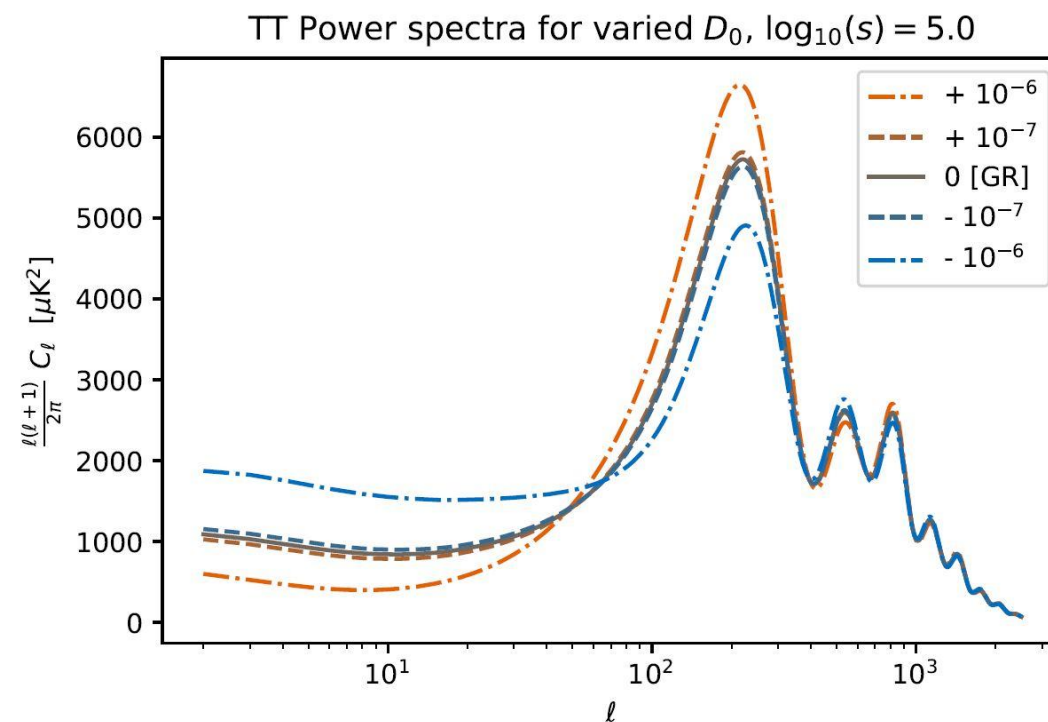
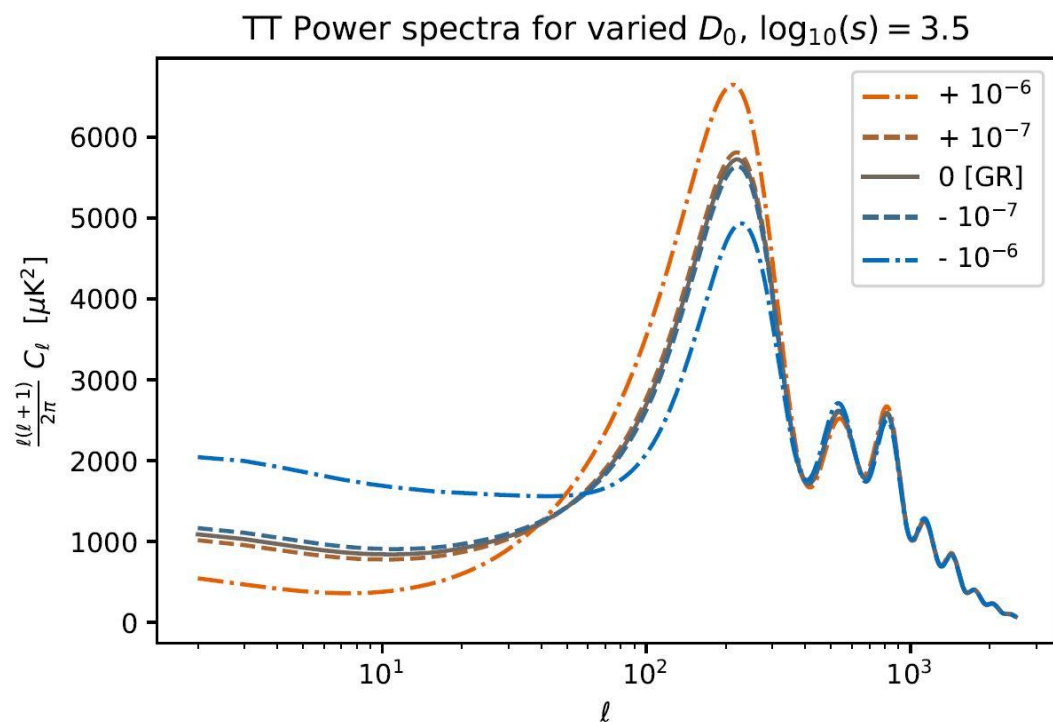
$$\Phi = (2D + 1)\Psi$$

- Difference between the potentials drops off as a Tanh function



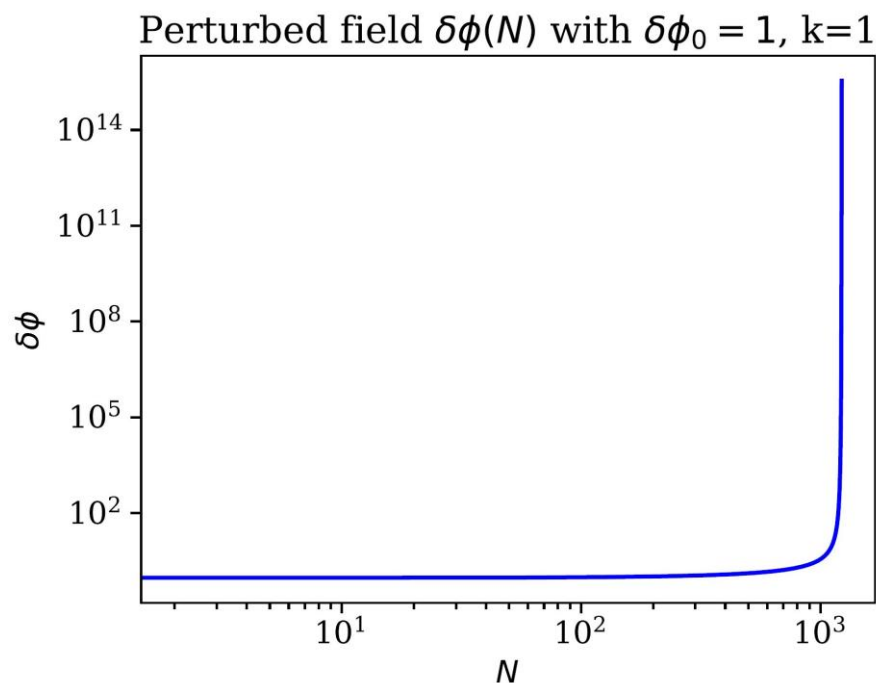
Previous New Higgs Results

- Affect on CMB for different s (rate of decay) values



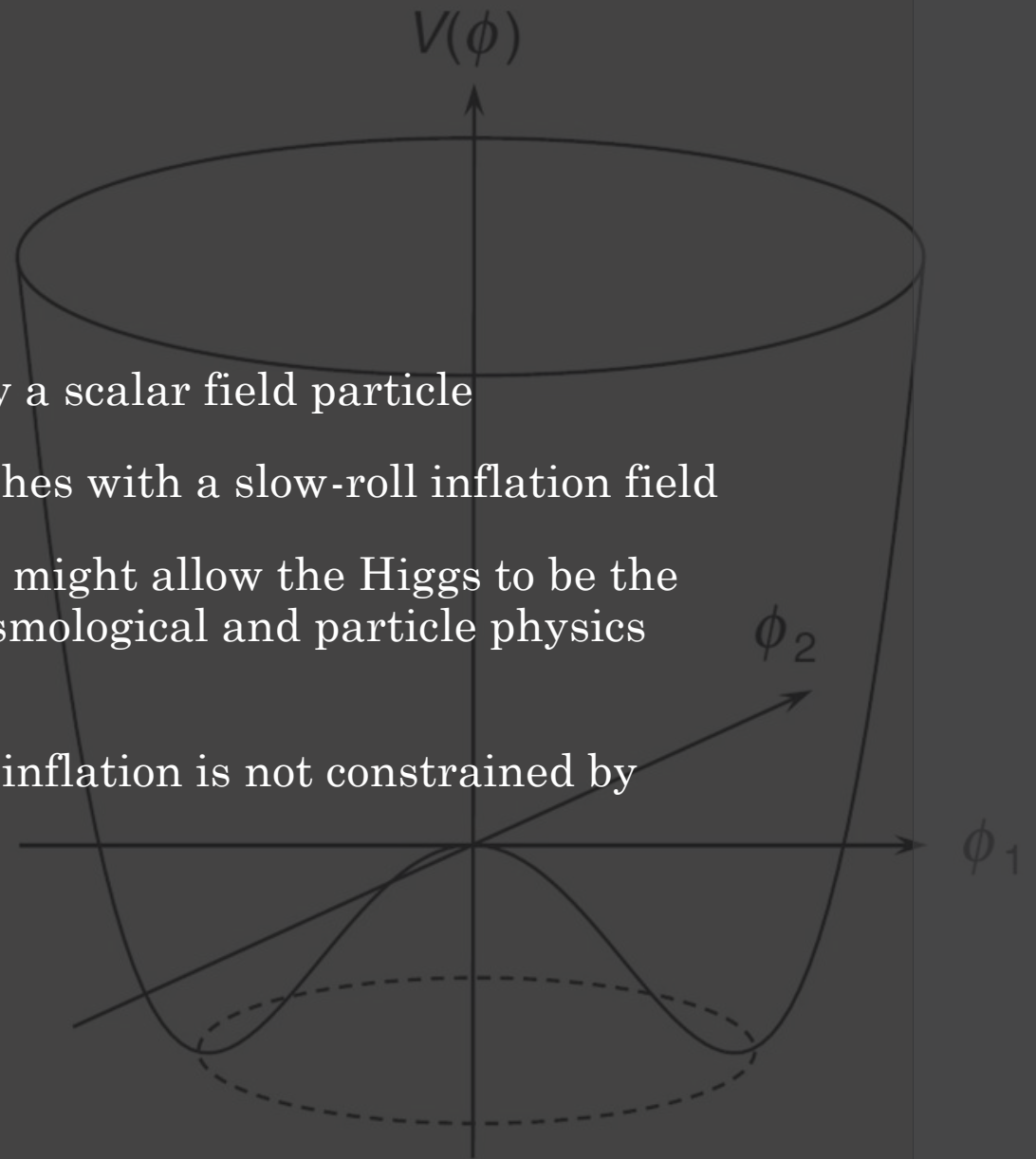
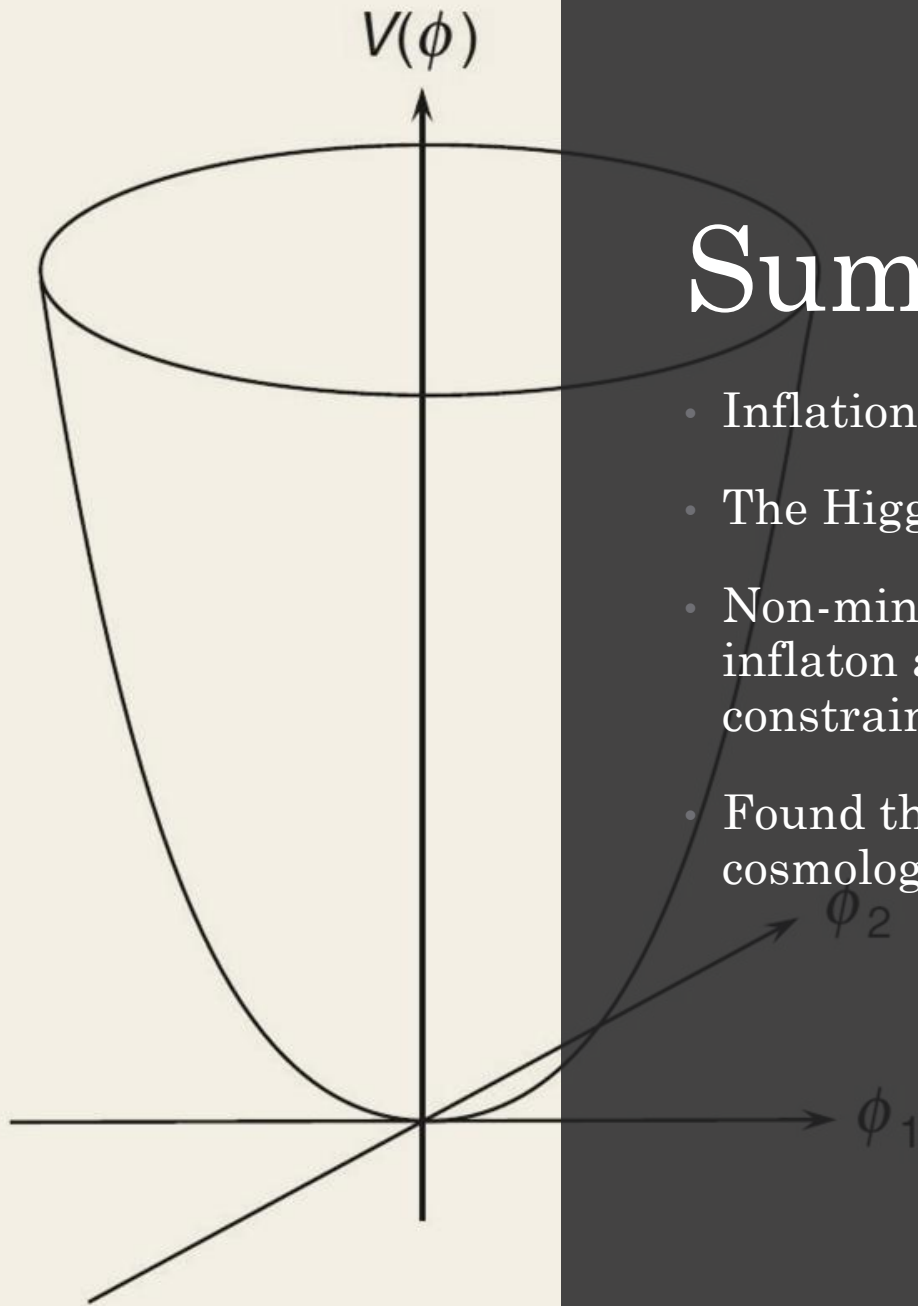
What Comes Next: Perturbations

- End goal is to find relationship between Φ and time-like Ψ during the early universe
- To do so, model perturbations around homogeneous equations
- Need to rigorously solve initial conditions for Old and New Higgs models



Summary

- Inflation is modeled by a scalar field particle
- The Higgs Boson matches with a slow-roll inflation field
- Non-minimal coupling might allow the Higgs to be the inflaton and match cosmological and particle physics constraints
- Found that Old Higgs inflation is not constrained by cosmology



The background of the slide is a grayscale image of the Cosmic Microwave Background (CMB) fluctuation map. It shows a curved, horizon-like section of the universe with a mottled texture of dark and light gray regions, representing temperature variations. The left side of the image is cut off by the edge of the frame.

Acknowledgements

- Thank you to my advisor Daniel Grin.
- David Kaiser for his Higgs Inflation notes
- Alan Guth for pioneering and developing this subject
- Our collaborators
- To Haverford KINSC office for funding me
- Questions?